

as $L \uparrow \Rightarrow m \uparrow$

$\omega_n = \sqrt{\frac{k}{m}}$ as $m \uparrow \omega_n \downarrow \Rightarrow$ larger cavity size $\Rightarrow$ lower resonant / natural frequency.

Finally in EM terms

~~$\omega_n = \sqrt{\frac{k}{L}}$~~ $\omega_n = \sqrt{LC}$

Journey of cylindrical cavity resonator

Pit stop Circular waveguide

cylindrical geometry $\Rightarrow$ cylindrical coordinates

$\phi \rightarrow$ angular component

$z \rightarrow$ axial component

$r \rightarrow$ radial component

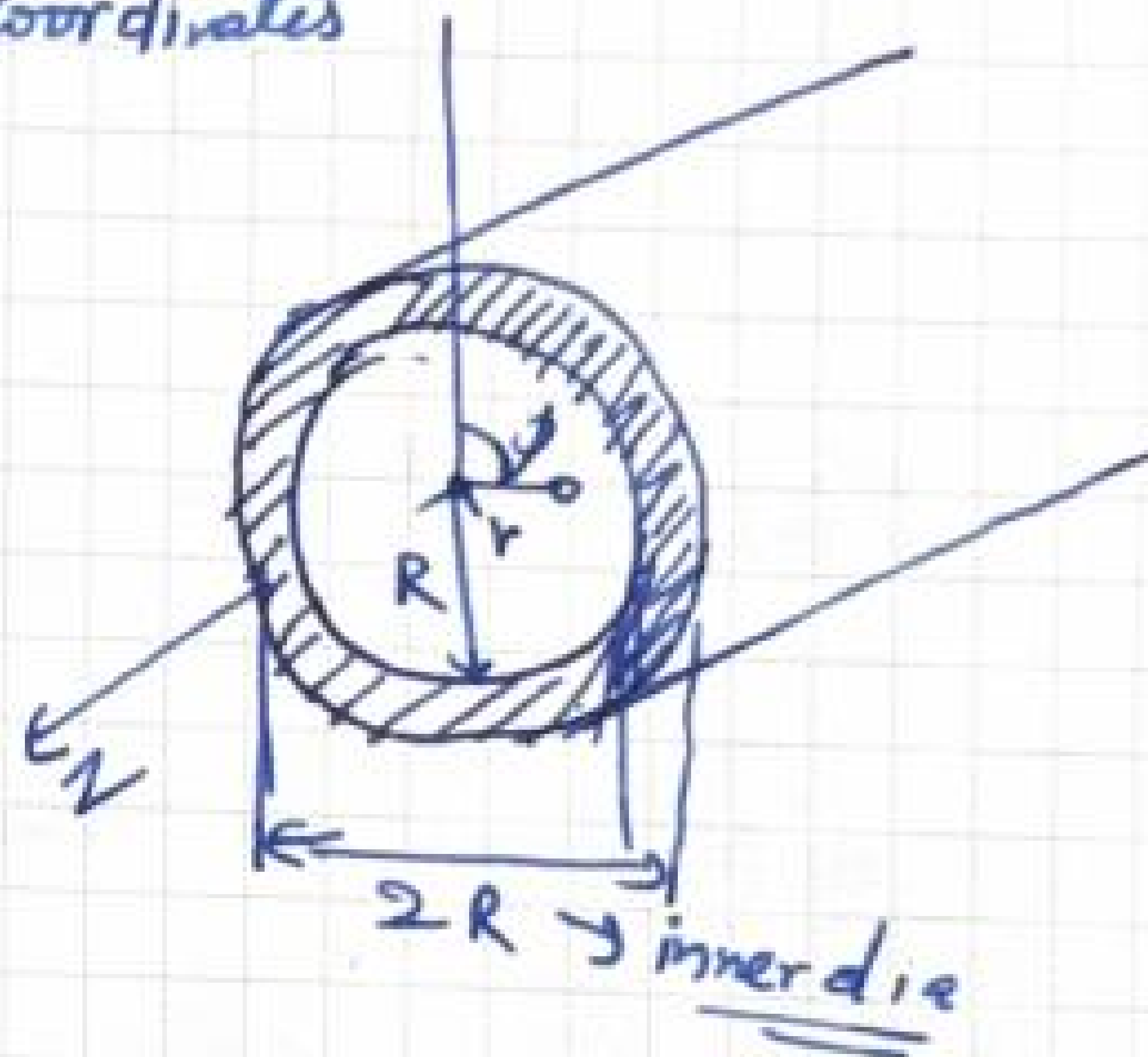
For dielec to conductor

Boundary $E_{\tan} = 0$
at surface

$$E_{\phi} = 0 @ r = R$$

$$E_z = 0 @ r = R$$

@ $r = R$ the boundary is conductor



For analyzing this we need to

Calculate $\vec{E}$ & $\vec{H}$ in this geometry

Maxwell Eqs + Boundary conditions will solve this geometry's $\vec{E}$ & $\vec{H}$

No 2 conductors, same conductor for return path $\Rightarrow$ TEM doesn't exist

Circular WG has TE & TM modes only

$z \nearrow$ $\phi \curvearrowright$ Both exist on surface of conductor at $r = R$

In WG $\nabla \cdot D = 0$ & $\nabla \cdot B = 0$ $\rightarrow j_v = 0 \rightarrow$ No bulk current

Source free $\Rightarrow$ divergence = 0

only 2 Maxwell eqn needed to be solved $\rightarrow$ coz this WG is a source free homogeneous region

$$\nabla \times E = -\mu \frac{dH}{dt}$$

$$\nabla \times H = \epsilon \frac{\partial E}{\partial t} + \bar{J}$$

Need to use this with cylindrical coordinates

$$\begin{aligned} \nabla \times E = & \begin{vmatrix} \frac{1}{r} \hat{r} & \hat{\phi} & \frac{1}{r} \hat{z} \\ \frac{\partial}{\partial r} & \frac{\partial}{\partial \phi} & \frac{\partial}{\partial z} \\ E_r & rE_{\phi} & E_z \end{vmatrix} = \frac{1}{r} \left[\frac{\partial E_z}{\partial \phi} - r \frac{\partial E_{\phi}}{\partial z} \right] \hat{r} \\ & - \left[\frac{\partial E_z}{\partial r} - \frac{\partial E_r}{\partial z} \right] \hat{\phi} \\ & + \frac{1}{r} \left[r \frac{\partial E_{\phi}}{\partial r} - \frac{\partial E_r}{\partial \phi} \right] \hat{z} \\ & = -\mu \frac{dH}{dt} \xrightarrow[\text{Fourier}]{\text{Laplace transform}} j\omega H \end{aligned}$$

Separating all vectors

$$\frac{1}{r} \frac{\partial E_z}{\partial \phi} - \frac{\partial E_{\phi}}{\partial z} = -\mu j\omega H_r \rightarrow \hat{r} \quad (1)$$

$$-\frac{\partial E_z}{\partial r} + \frac{\partial E_r}{\partial z} = -\mu j\omega H_{\phi} \rightarrow \hat{\phi} \quad (2)$$

$$\frac{\partial E_{\phi}}{\partial r} - \frac{1}{r} \frac{\partial E_r}{\partial \phi} = -\mu j\omega H_z \rightarrow \hat{z} \quad (3)$$

$$\nabla \times H = \begin{vmatrix} \frac{1}{r} \hat{r} & \hat{\phi} & \frac{1}{r} \hat{z} \\ \frac{\partial}{\partial r} & \frac{\partial}{\partial \phi} & \frac{\partial}{\partial z} \\ H_r & r H_\phi & H_z \end{vmatrix} = \frac{1}{r} \left[\frac{\partial H_z}{\partial \phi} - r \frac{\partial H_\phi}{\partial z} \right] \hat{r} - \left[\frac{\partial H_z}{\partial r} - \frac{\partial H_r}{\partial z} \right] \hat{\phi} + \frac{1}{r} \left[r \frac{\partial H_\phi}{\partial r} - \frac{\partial H_r}{\partial \phi} \right] \hat{z}$$

$$= \epsilon \left(\frac{\partial \vec{E}}{\partial t} \right) + \vec{J}$$

Fourier transform $\rightarrow$ $\vec{J} = \sigma \vec{E}$
 $J = 0$ Since wave is hollow $\Rightarrow \sigma \rightarrow 0$ inside

Separating all vectors

$$\frac{1}{r} \frac{\partial H_z}{\partial \phi} - \frac{\partial H_\phi}{\partial z} = \epsilon j \omega E_r \rightarrow \hat{r} \quad (4)$$

$$-\frac{\partial H_z}{\partial r} + \frac{\partial H_r}{\partial z} = \epsilon j \omega E_\phi \rightarrow \hat{\phi} \quad (5)$$

$$\frac{\partial H_\phi}{\partial r} - \frac{1}{r} \frac{\partial H_r}{\partial \phi} = \epsilon j \omega E_z \rightarrow \hat{z} \quad (6)$$

observing 6 equations — The solution must define a propagating wave

The wave (EM wave) in the waveguide travels along 'z' direction.

The general phasor form for travelling wave is — travelling in 'z' direction

$$\vec{E}(r, \phi, z) = \vec{E}(r, \phi) e^{-j\gamma z}$$

$$H(r, \phi, z) = H(r, \phi) e^{-j\gamma z}$$

$$\gamma = \alpha + j\beta$$

$\downarrow$
if lossless medium
 $\alpha = 0$

Assume lossless medium

$$\vec{E}(r, \phi, z) = \vec{E}(r, \phi) e^{-j\beta z}$$

$$H(r, \phi, z) = \vec{H}(r, \phi) e^{-j\beta z}$$

Both $\vec{E}$ & $\vec{H}$ are functions of r, ϕ & propagate in 'z'

$\therefore \nabla \times \vec{E}$

Comparing differential & phasor form

we know that

$$\frac{\partial}{\partial z} \rightarrow -jk$$

$$\left(\frac{\partial}{\partial z} e^{-jkz} = -jk e^{-jkz} \right)$$

Interchange const

$$\therefore \frac{\partial}{\partial z} \rightarrow e^{-jkz} -jk$$

'- represents travelling in forward direction

Now lets replace $\frac{\partial}{\partial z}$ by $-j\beta$

Take eq 2 & 4 pair (both are orthogonal in direction)

$$-\frac{\partial E_z}{\partial r} + \frac{\partial E_r}{\partial z} = -\mu j\omega H_\phi \rightarrow -\frac{\partial E_z}{\partial r} - j\beta E_r = -\mu j\omega H_\phi \quad (7)$$

$$\frac{1}{r} \frac{\partial}{\partial \phi} H_z + j\beta H_\phi = \epsilon j\omega E_r \quad (8)$$

lets find 'E' fields first $\Rightarrow$ replace to eliminate 'H' terms

$$H_\phi = -\frac{1}{\mu j\omega} \left[-\frac{\partial E_z}{\partial r} + \frac{\partial E_r}{\partial z} \right] \times \frac{1}{-j\omega\mu}$$

$$H_\phi = \frac{j}{\mu\omega} \left[\frac{\partial E_z}{\partial r} - \frac{\partial E_r}{\partial z} \right] \rightarrow \text{Substitute this in}$$

Confused & wrong!

$$\frac{1}{r} \frac{\partial}{\partial \phi} H_z + j\beta \left[\frac{j}{\mu\omega} \left(\frac{\partial E_z}{\partial r} - \frac{\partial E_r}{\partial z} \right) \right] = \epsilon j\omega E_r$$

$$\frac{1}{r} \frac{\partial}{\partial \phi} H_z + \frac{j\beta}{\mu\omega} \left[\frac{\partial E_z}{\partial r} - \frac{\partial E_r}{\partial z} \right] = \epsilon j\omega E_r$$

$$\frac{1}{\epsilon j\omega r} \frac{\partial}{\partial \phi} H_z + \frac{j\beta}{\mu\omega \epsilon j\omega} \left[\frac{\partial E_z}{\partial r} - \frac{\partial E_r}{\partial z} \right] = E_r$$

$$\frac{-j}{\epsilon\omega r} \frac{\partial}{\partial \phi} H_z + \frac{\beta}{\omega^2 \mu \epsilon} \left[\frac{\partial E_z}{\partial r} - \frac{\partial E_r}{\partial z} \right] = E_r$$

off by 1

$$H_\phi = \frac{1}{\mu j\omega} \frac{\partial}{\partial r} E_z - \frac{j\beta E_r}{-\mu j\omega} \quad \text{--- substitute this in eqn ②}$$

$$H_\phi = \frac{-j}{\omega\mu} \frac{\partial}{\partial r} E_z + \beta \frac{E_r}{\mu\omega}$$

$$\Rightarrow \frac{1}{r} \frac{\partial}{\partial \phi} H_z + j\beta \left[\frac{-j}{\omega\mu} \frac{\partial}{\partial r} E_z + \frac{\beta E_r}{\omega\mu} \right] = \epsilon j\omega E_r$$

$$\frac{1}{r\epsilon j\omega} \frac{\partial}{\partial \phi} H_z + \frac{j\beta}{j\omega\epsilon} \times \frac{1}{\omega\mu} \left[-j \frac{\partial}{\partial r} E_z + \beta E_r \right] = E_r$$

$$\frac{1}{r\epsilon j\omega} \frac{\partial}{\partial \phi} H_z + \frac{j\beta}{\omega^2\mu\epsilon} \frac{\partial}{\partial r} E_z + \frac{\beta^2}{\mu\epsilon\omega^2} E_r = E_r$$

$$\frac{1}{r\epsilon j\omega} \frac{\partial}{\partial \phi} H_z - \frac{j\beta}{\omega^2\mu\epsilon} \frac{\partial}{\partial r} E_z = \left[1 - \frac{\beta^2}{\omega^2\mu\epsilon} \right] E_r$$

$$E_r = \frac{1}{r\epsilon j\omega} \left[1 - \frac{\beta^2}{\omega^2\mu\epsilon} \right] \times \frac{\partial}{\partial \phi} H_z - \frac{j\beta}{\omega^2\mu\epsilon \left[1 - \frac{\beta^2}{\omega^2\mu\epsilon} \right]} \times \frac{\partial}{\partial r} E_z$$

$$E_r = \frac{-j\omega\mu}{r\omega\epsilon \left[\omega^2\mu\epsilon - \beta^2 \right]} \times \frac{\partial H_z}{\partial \phi} - \frac{j\beta}{\omega^2\mu\epsilon - \beta^2} \times \frac{\partial E_z}{\partial r}$$

$$E_r = \frac{-j\omega\mu}{r \left[\omega^2\mu\epsilon - \beta^2 \right]} \frac{\partial H_z}{\partial \phi} - \frac{j\beta}{\omega^2\mu\epsilon - \beta^2} \frac{\partial E_z}{\partial r}$$

$\swarrow$ [let $k_c^2 = \omega^2\mu\epsilon - \beta^2 \rightarrow$ this is called wave number cut off
off track let's go into units of this

$\beta \rightarrow$ phase constant $\rightarrow$ longitudinal wave number
 How fast phase changes over distance
 units rad/m $\rightarrow$ there is z

$\epsilon C = Q/V \quad I = \frac{Q}{t}$
 Coulomb x Vol
 Amp x Sec x V
 $V = L \frac{dI}{dt} \quad \frac{V \times \text{Sec}}{\text{Amp}}$

$\beta^2 \rightarrow \left(\frac{\text{rad}}{\text{m}}\right)^2$

$\mu \rightarrow \frac{\text{Farad}}{\text{m}} \quad \text{Henry}$
 $\frac{\text{Volt} \times \text{Sec}}{\text{Amp}}$

$\epsilon = \frac{\text{Farad}}{\text{m}} \rightarrow \frac{\text{Amp} \times \text{Sec}}{\text{Volt}}$

$\omega^2 \rightarrow \frac{\text{rad}^2}{\text{Sec}^2}$

$\rightarrow$ Add / Sub unit should be equal

$\beta^2 \rightarrow \left(\frac{\text{rad}}{\text{m}}\right)^2 = \frac{\text{rad}^2}{\text{Sec}^2} \times \frac{\text{Henry}}{\text{m}} \times \frac{\text{Farad}}{\text{m}}$
 $\rightarrow \text{WHE}$

$\frac{\text{rad}^2}{\text{Sec}^2} \times \frac{\text{V} \times \text{Sec}}{\text{Amp} \times \text{m}} \times \frac{\text{Amp} \times \text{Sec}}{\text{Vol} \times \text{m}}$

$\left(\frac{\text{rad}}{\text{m}}\right)^2 = \frac{\text{rad}^2}{\text{m}^2}$ unit match

k_c^2 units $\frac{\text{rad}^2}{\text{m}^2} \rightarrow$ Cut off wave number
 (Transverse wave number)

that's why k_c^2 we assumed not ' k_c '

$k_c \rightarrow$ termed for Transverse wave number / cut off wave number

gold bar
 k in WB $\rightarrow$

$$E_r = \frac{-j\omega\mu}{rk_c^2} \frac{\partial H_z}{\partial \phi} - \frac{j\beta}{k_c^2} \frac{\partial E_z}{\partial r}$$

Now using Eqn 14.5 for E_ϕ component

1 $\rightarrow \frac{1}{r} \frac{\partial E_z}{\partial \phi} - (-j\beta) E_\phi = -\mu j\omega H_r \quad \text{--- (9)}$

5 $\rightarrow -\frac{\partial H_z}{\partial r} + (-j\beta) H_r = \epsilon j\omega E_\phi \quad \text{--- (10)}$

From ①

$$H_r = \frac{-1}{rj\mu\omega} \frac{\partial E_z}{\partial \phi} - \frac{j\beta}{-j\mu\omega} E\phi$$

$$H_r = \frac{j}{r\mu\omega} \frac{\partial E_z}{\partial \phi} - \frac{\beta}{\mu\omega} E\phi$$

↳ substitute this in eqn 5

$$-\frac{\partial H_z}{\partial r} - j\beta \left[\frac{j}{r\mu\omega} \frac{\partial E_z}{\partial \phi} - \frac{\beta}{\mu\omega} E\phi \right] = \epsilon j\omega E\phi$$

$$-\frac{\partial H_z}{\partial r} + \frac{j\beta}{r\mu\omega} \frac{\partial E_z}{\partial \phi} + \frac{j\beta^2}{\mu\omega} E\phi = \epsilon j\omega E\phi$$

$$E\phi \left[\epsilon j\omega + \frac{j\beta^2}{\mu\omega} \right]$$

$$E\phi = -\frac{\partial H_z}{\partial r} \left[\frac{1}{j[\epsilon\omega + \frac{\beta^2}{\mu\omega}]} \right] + \frac{j\beta}{j r \mu \omega [\epsilon\omega + \frac{\beta^2}{\mu\omega}]} \frac{\partial E_z}{\partial \phi}$$

$$\frac{\partial H_z}{\partial r} \left[\frac{j\mu\omega}{\mu\epsilon\omega^2 + \beta^2} \right] - \frac{j\beta}{r\mu\omega^2 + r\beta^2} \frac{\partial E_z}{\partial \phi}$$

$$E\phi = \frac{\partial H_z}{\partial r} \left[\frac{j\omega\mu}{\mu\epsilon\omega^2 - \beta^2} \right] - \frac{j\beta}{r[\mu\omega^2 - \beta^2]} \frac{\partial E_z}{\partial \phi}$$

$$\text{again } k_c^2 = \omega^2\mu\epsilon - \beta^2$$

$$E\phi = \frac{j\omega\mu}{k_c^2} \frac{\partial H_z}{\partial r} - \frac{j\beta}{r k_c^2} \frac{\partial E_z}{\partial \phi}$$

$$E_\phi = -\frac{j}{k_c^2} \left[\frac{\beta}{r} \frac{\partial E_z}{\partial \phi} - \omega \mu \frac{\partial H_z}{\partial r} \right]$$

H-fields derivation

Let's derive 'H_φ'

$$7 - \quad -\frac{\partial}{\partial r} E_z - j\beta E_r = -\mu j\omega H_\phi$$

$$8 \rightarrow \quad \frac{1}{r} \frac{\partial}{\partial \phi} H_z + j\beta H_\phi = \epsilon j\omega E_r$$

$$-\frac{\partial}{\partial r} E_z - j\beta \left[\frac{1}{r\epsilon j\omega} \frac{\partial}{\partial \phi} H_z + \frac{j\beta}{\epsilon j\omega} H_\phi \right] = -\mu j\omega H_\phi$$

$$-\frac{\partial}{\partial r} E_z - \frac{j\beta}{r\epsilon j\omega} \frac{\partial}{\partial \phi} H_z + \frac{j\beta^2}{\epsilon j\omega} H_\phi = -\mu j\omega H_\phi$$

$$\left(-\mu j\omega + \frac{j\beta^2}{\epsilon j\omega} \right) H_\phi$$

$$\left[-\frac{\partial}{\partial r} E_z - \frac{\beta}{r\epsilon j\omega} \frac{\partial}{\partial \phi} H_z \right] = H_\phi \left[\frac{j\beta^2}{\epsilon j\omega} - j\omega\mu \right]$$

$$\left(\frac{-1}{\frac{j\beta^2}{\epsilon j\omega} - j\omega\mu} \right) \frac{\partial}{\partial r} E_z - \frac{\beta}{r\epsilon j\omega} \left[\frac{j\beta^2}{\epsilon j\omega} - j\omega\mu \right] \frac{\partial H_z}{\partial \phi} = H_\phi$$

$$\frac{-1}{j \left(\frac{\epsilon\omega}{\beta^2 - \omega^2\mu\epsilon} \right)} \frac{\partial}{\partial r} E_z - \frac{-\beta}{j(\beta^2 - \omega^2\mu\epsilon r)} \frac{\partial H_z}{\partial \phi} = H_\phi$$

$$\frac{j\beta}{r k_c^2} \frac{\partial H_z}{\partial \phi} - \frac{j\epsilon\omega}{k_c^2} \frac{\partial E_z}{\partial r} = H_\phi$$

$$\boxed{-\frac{j}{kc^2} \left[\frac{\beta}{r} \frac{\partial H_z}{\partial \phi} + \epsilon \omega \frac{\partial E_z}{\partial r} \right] = H_\phi}$$

'Hr' derive

$$\frac{1}{r} \frac{\partial E_z}{\partial \phi} - (-j\beta) E_\phi = -\mu j\omega H_r \quad - (9)$$

$$-\frac{\partial H_z}{\partial r} + (-j\beta) H_r = \epsilon j\omega E_\phi \quad - (10)$$

From (9) $H_r = \frac{1}{-r j\omega \mu} \frac{\partial E_z}{\partial \phi} + \frac{j\beta}{-j\omega \mu} E_\phi$

$H_r = \frac{+j}{\omega \mu r} \frac{\partial E_z}{\partial \phi} - \frac{\beta}{\omega \mu} E_\phi$ Confused again!
 Substitute this in (10) Redial E_ϕ

$$-\frac{\partial H_z}{\partial r} + (-j\beta) \left[\frac{j}{\omega \mu r} \frac{\partial E_z}{\partial \phi} - \frac{\beta}{\omega \mu} E_\phi \right] = \epsilon j\omega E_\phi$$

$$-\frac{\partial H_z}{\partial r} - \left[\frac{-\beta}{\omega \mu r} \frac{\partial E_z}{\partial \phi} \right] - \left[\frac{-j\beta^2}{\omega \mu} E_\phi \right] = \epsilon j\omega E_\phi$$

$$-\frac{\partial H_z}{\partial r} + \frac{\beta}{\omega \mu r} \frac{\partial E_z}{\partial \phi} = E_\phi \left[\epsilon j\omega - \frac{j\beta^2}{\omega \mu} \right]$$

$$E_\phi = \left(\frac{-1}{\left(\epsilon j\omega - \frac{j\beta^2}{\omega \mu} \right)} + \frac{\beta}{\omega \mu r \left[\epsilon j\omega - \frac{j\beta^2}{\omega \mu} \right]} \right) \frac{\partial E_z}{\partial \phi}$$

$$-\frac{\partial H_z}{\partial r} \left(\frac{1}{\epsilon j \omega} \right) + \frac{(-j\beta)}{\epsilon j \omega} H_r = E \phi$$

$$-\frac{\partial H_z}{\partial r} \left(\frac{-j}{\epsilon \omega} \right) - \frac{\beta}{\epsilon \omega} H_r = E \phi$$

Sub this in 9

$$\frac{1}{r} \frac{\partial E_z}{\partial \phi} + j\beta \left[\frac{j}{\epsilon \omega} \frac{\partial H_z}{\partial r} - \frac{\beta}{\epsilon \omega} H_r \right] = -j\omega \mu H_r$$

$$\frac{1}{r} \frac{\partial E_z}{\partial \phi} - \frac{\beta}{\epsilon \omega} \frac{\partial H_z}{\partial r} - \frac{j\beta^2}{\epsilon \omega} H_r = -j\omega \mu H_r$$

$$H_r \left[\frac{j\beta^2}{\epsilon \omega} - j\omega \mu \right]$$

$$\frac{1}{r \left[\frac{j\beta^2}{\epsilon \omega} - j\omega \mu \right]} \frac{\partial E_z}{\partial \phi} - \frac{\beta}{\epsilon \omega \left[\frac{j\beta^2}{\epsilon \omega} - j\omega \mu \right]} \frac{\partial H_z}{\partial r} = H_r$$

$$\frac{\epsilon \omega}{j r [\beta^2 - \omega^2 \mu \epsilon]} \frac{\partial E_z}{\partial \phi} - \frac{\beta}{j [\beta^2 - \omega^2 \mu \epsilon]} \frac{\partial H_z}{\partial r} = H_r$$

$$\frac{j \epsilon \omega}{r k_c^2} \frac{\partial E_z}{\partial \phi} - \frac{j \beta}{k_c^2} \frac{\partial H_z}{\partial r} = H_r$$

$$H_r = \frac{j}{k_c^2} \left[\frac{\epsilon \omega}{r} \frac{\partial E_z}{\partial \phi} - \beta \frac{\partial H_z}{\partial r} \right]$$

Writing All 4 eqns E_r, E_ϕ, H_r, H_ϕ

$$E_r = -\frac{j}{k_c^2} \left[\frac{\omega\mu}{r} \frac{\partial H_z}{\partial \phi} + \beta \frac{\partial E_z}{\partial r} \right]$$

$$E_\phi = -\frac{j}{k_c^2} \left[\frac{\beta}{r} \frac{\partial E_z}{\partial \phi} - \omega\mu \frac{\partial H_z}{\partial r} \right]$$

$$H_\phi = -\frac{j}{k_c^2} \left[\frac{\beta}{r} \frac{\partial H_z}{\partial \phi} + \epsilon\omega \frac{\partial E_z}{\partial r} \right]$$

$$H_r = -\frac{j}{k_c^2} \left[\beta \frac{\partial H_z}{\partial r} - \frac{\epsilon\omega}{r} \frac{\partial E_z}{\partial \phi} \right]$$

What about E_z & H_z values?

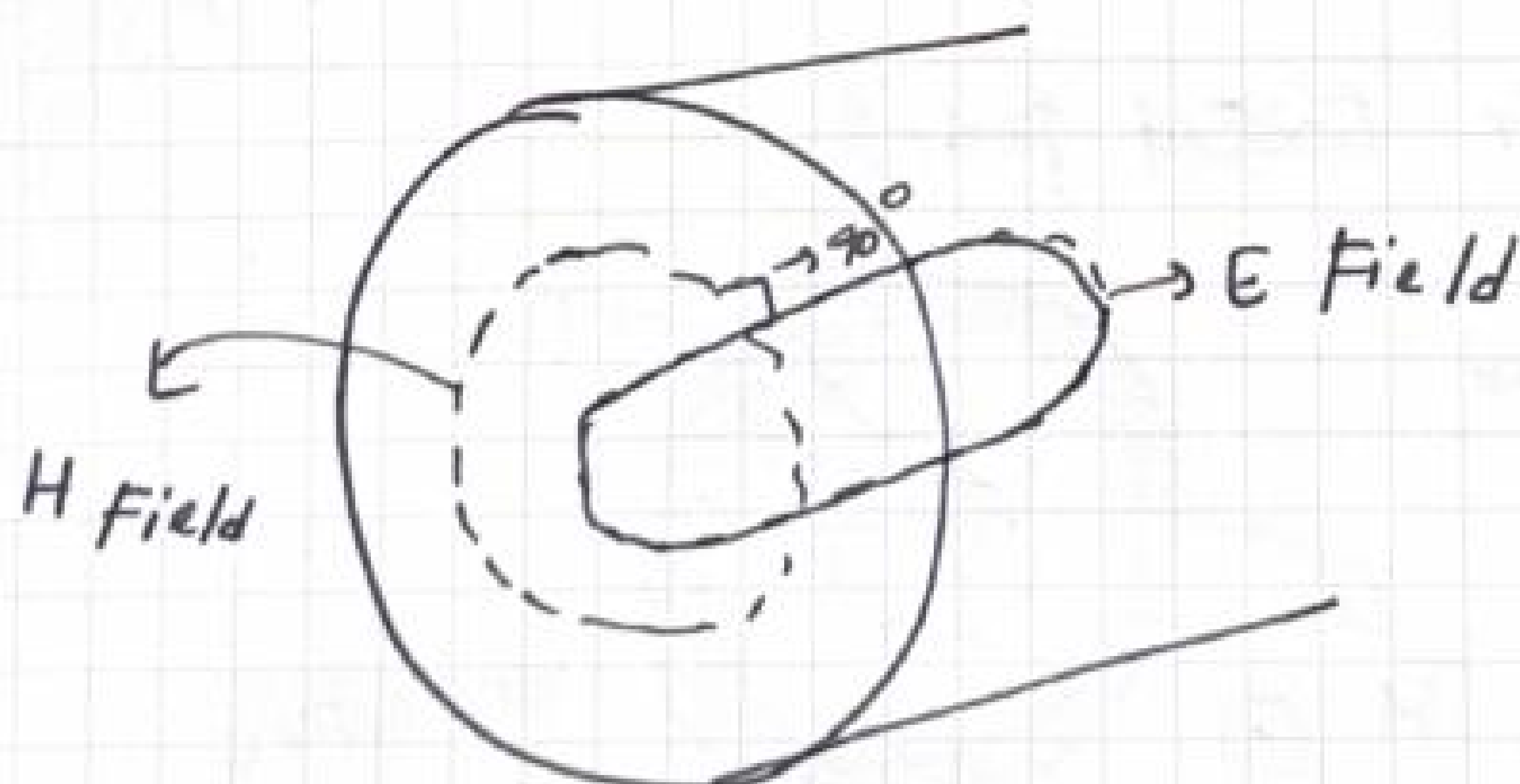
First of all let's observe one thing, how the wave propagates.

If it has to propagate in TEM then both E_z & $H_z = 0$

if this happens, the above 4 eqns $E_r, E_\phi, H_r, H_\phi \rightarrow 0$

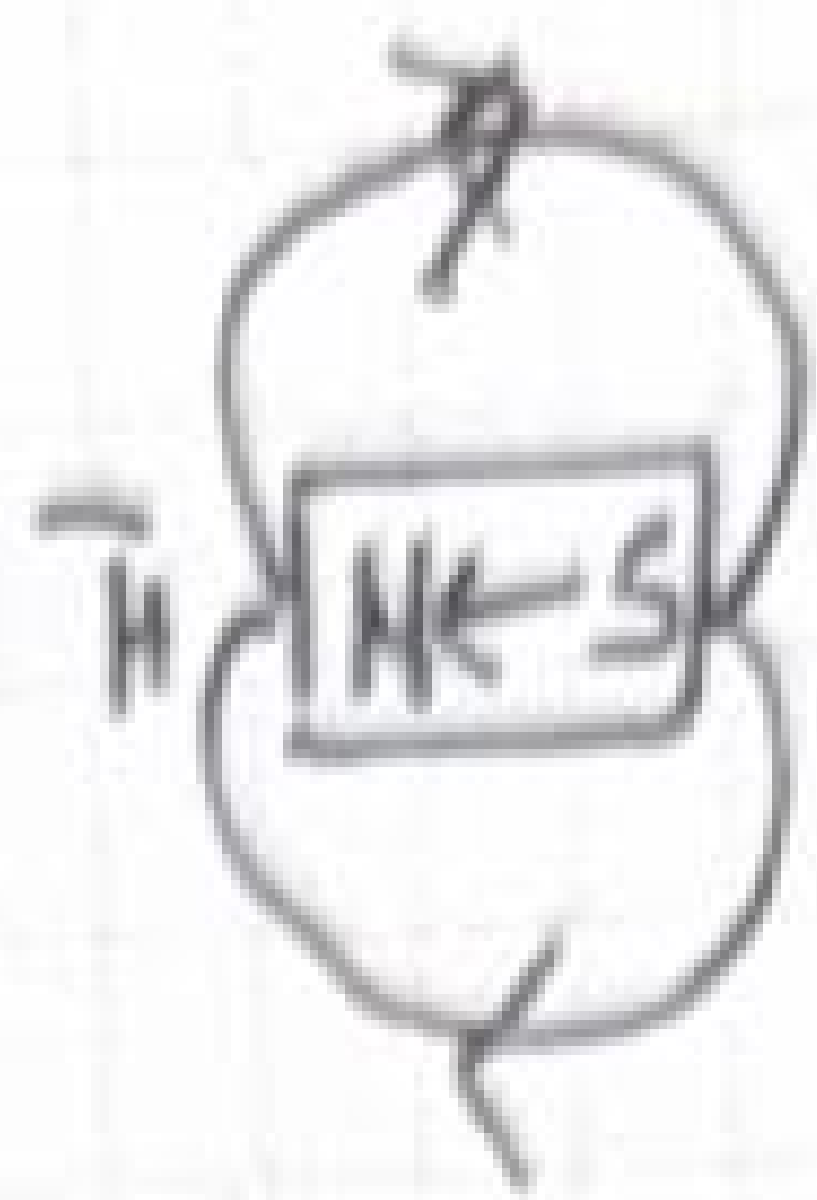
So, TEM doesn't exist - Mathematically proved!

Let's understand this physically!



Rule
 $\vec{E} \perp \vec{H}$
 always

We know from fundamentals $\vec{H}$ is always in loop



$\vec{E}$ is from +ve charge to -ve charge $+Q \xrightarrow{\vec{E}} -Q$

If $\vec{H}$ is looping around the cross section of wg

then $\vec{E}$ has to be $\perp$ to $\vec{H}$

in a hollow metal tube there is no potential difference

so $\vec{E}$ has to curl & form loop no choice

The only possible $\perp$ loop $\vec{E}$ can form has to be $\perp$

Cross section $\Rightarrow$ definitely it has $\vec{z}$ component

$z \rightarrow$ direction of wave propagation.

So, Longitudinal component exists always.

Same way if $\vec{E}$ curls around cross section $\vec{H}$

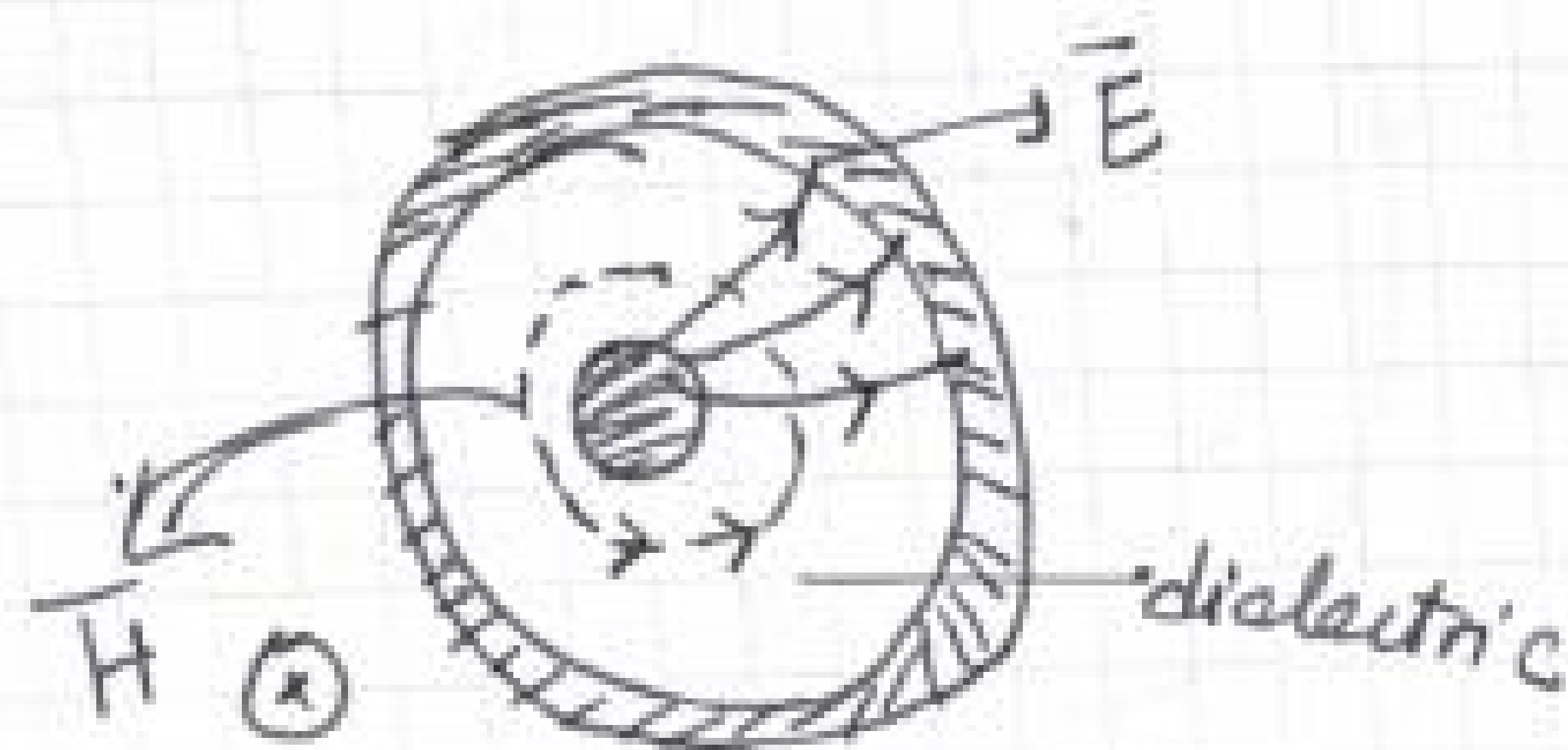
has no choice to curl $\perp$ to it & composing of longitudinal component.

So, fully TEM cannot exist in hollow metallic conductor AKA Waveguide:

Why TEM can exist in a 2 conductor system

Like coax or twisted pair?

Lets take coax



In a coax the inner & outer conductors are separated by dielectric.

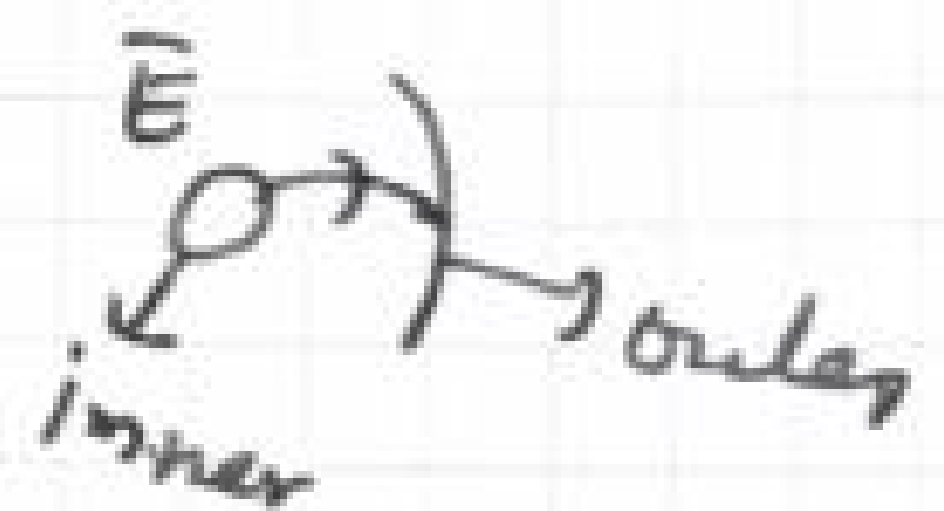
The inner conductor is at a different potential compared to outer conductor forming a potential difference.

When there is potential difference, charges in the conductor flow forming a magnetic loop around center conductor. (No Hz component)

Since there is potential difference b/w the conductors, the charges try to attract from one conductor to another forming an $\vec{E}$ field radiating radially out from center to outer conductor. (~~displacement current~~).

$\vec{E}$ doesn't loop in itself but loops $\rightarrow$ no loop is axially formed. is one direction

$\therefore$ there is no longitudinal component
 $\vec{E}_z$ & H_z .



$\therefore$ TEM wave exists in 2 conductor system!